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MEASURING COLLIMATOR INFRARED (IR) SPECTRAL TRANSMISSION

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and Engineering Center

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TABLE OF CONTENTS

| | <u>Page</u> |
|---|-------------|
| I. INTRODUCTION | 1 |
| II. BACKGROUND..... | 1 |
| III. SETUP AND TESTING..... | 2 |
| IV. RESULTS AND ANALYSIS | 7 |
| V. CONCLUSION AND FUTURE WORK..... | 8 |
| REFERENCES | 9 |
| LIST OF ABBREVIATIONS, ACRONYMS, AND SYMBOLS | 10 |

LIST OF ILLUSTRATIONS

| <u>Figure</u> | <u>Title</u> | <u>Page</u> |
|----------------------|--|--------------------|
| 1. | Collimated, Distance, and Flooded Test Configurations | 2 |
| 2. | Collimator Transmission | 3 |
| 3. | Setup of SR5000 With Atmospheric Paths Equal at 4 m..... | 4 |
| 4. | Response Function at 100 °C | 5 |
| 5. | Response at Distance Divided by Response at Infinity | 6 |
| 6. | Change in F/# Due to Change in Focus for Each Distance | 6 |
| 7. | Radiance at 100 °C | 7 |
| 8. | Radiance at 100 °C Zoomed-In | 7 |
| 9. | Reported Versus Measured Collimator Transmission..... | 8 |

LIST OF TABLES

| <u>Table</u> | <u>Title</u> | <u>Page</u> |
|---------------------|--|--------------------|
| 1. | Parameters for Calculation Maximum Working Distance..... | 3 |
| 2. | Average Transmission of Several Sub-Bands | 8 |

I. INTRODUCTION

Several Infrared (IR) imaging systems have been measured by looking down a collimator and a standalone large surface blackbody. The total atmospheric path length between source and sensor were the same, yet discrepancies were noticed in that the standalone blackbody measurements were providing more perceived flux even when compensating for vendor-provided collimator transmission losses. This report provides the background and discusses the methodology used to measure the actual transmission losses in a collimator.

II. BACKGROUND

Several IR systems [1] have been measured looking down a collimator using a large surface blackbody with a single fold/relay mirror and an off axis parabolic mirror to simulate an object at infinity. These systems were also measured apart from the collimator by directly viewing a large surface blackbody up close. Both methods flood-filled the pixels of interest on the focal plane and had the blackbodies at identical temperatures. Discrepancies were noticed. The up-close blackbody measurements provided significantly more perceived flux at the detector level even when compensating for atmospheric path effects and vendor-provided transmission losses of the mirrors in the collimator. Mirrors actually do not transmit but rather reflect. In the context of this reflective collimator system, this is known as a transmission loss. Discrepancies were also noted when the large surface blackbody was placed at various distances from the sensors under test. At each distance the sensor was refocused and atmospheric losses were applied.

Most laboratories perform only one technique (Figure 1) for any given system, but rarely (if ever) measure the same system with multiple test setup configurations. This is mostly due to lack of available time and/or equipment. To eliminate the test equipment as the root cause, the first step was to repeat the measurement with a completely different standalone blackbody and collimator system. Similar results were achieved with the other blackbodies and collimators. To add to the data set, additional sensors were also evaluated to help determine potential trends and help point to a root cause. When these methods failed to point to an obvious cause, a more detailed drill down was initiated. The community literature was consulted and several lines of thought were discussed that could be causing the effect: atmospheric transmission calculations were not adequate, spectral emissivity differences of the two blackbodies, and larger than reported transmission losses of the collimator optics.

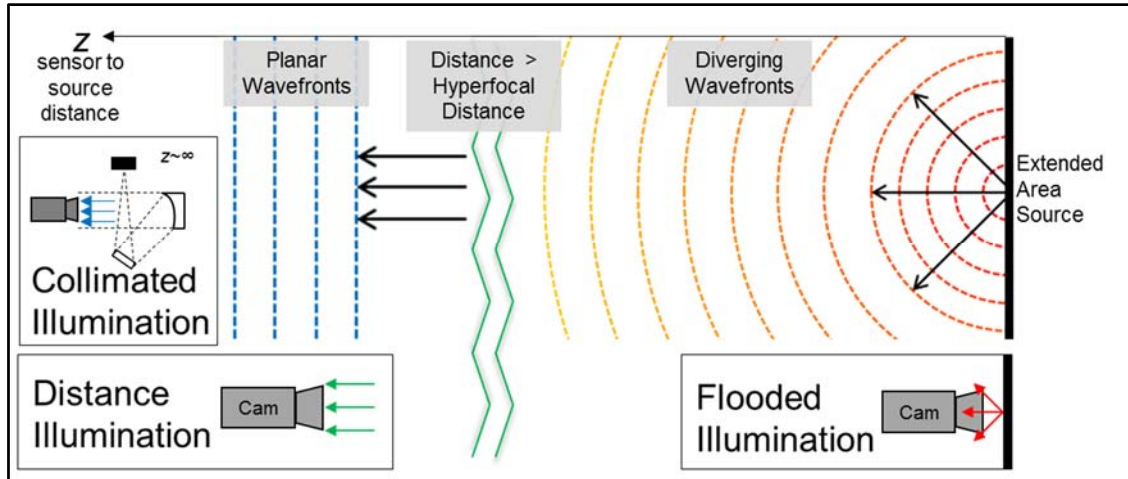


Figure 1. Collimated, Distance, and Flooded Test Configurations

III. SETUP AND TESTING

The collimator was investigated to its possible contributions to the measured discrepancies. When a separate collimator was giving similar results, a literature search was completed concerning the possible pitfalls of a reflective parabolic collimator. While a parabolic reflective system has many benefits, such as the ability to work over a broad spectral band (for example, 0.3 to 15 micrometers (μm)), no spherical aberrations, no chromatic aberrations, and large apertures, it does have a drawback if aluminum mirrors with protective silicon monoxide (SiO) coatings are used. Most transmission measurements of optics (in this case, mirrors) are reported at normal incidence angles, yet the fold mirror is used at an angle. Potential large losses [2,3,4] (approximately 15 percent (%)) are possible at approximately 8 μm with additional losses to nearly 10 μm before flattening out and returning to a normal 2 to 3% loss. This was the item to be investigated.

After consulting the collimator documentation, it was determined that the primary mirror had an aluminum coating with SiO as a protective layer. It was reported to have more than a 97.0% transmission from 3 to 15 μm . A scanned graph of a spectral curve was also provided, but detailed tabular data was lacking. The secondary mirror was also similar in its coating and documentation. Figure 2 shows a digitized version of the provided scanned graphs. ImageJ and Excel were used to digitize the scans and calculate the overall collimator system transmission.

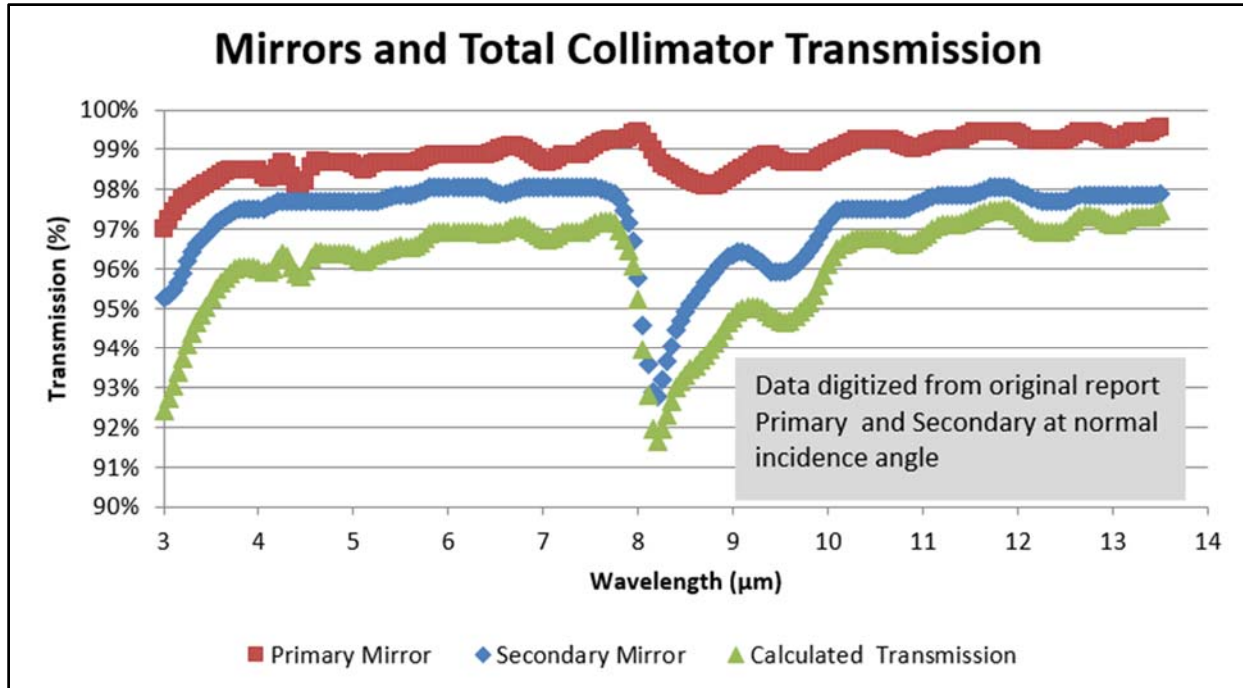


Figure 2. Collimator Transmission

Figure 2 shows that there is a large dip in the transmission at approximately 8 μm for the secondary mirror. This was at a normal incidence angle, and additional losses were expected based on the literature search and the fold mirror being at an angle. A spectral radiometer was then secured and used to measure the true spectral transmission. A CI Systems SR5000W with an indium antimonide (InSb)/mercury cadmium telluride (MCT) detector and continuous variable filter was selected as the instrument best suited to perform the measurement. The maximum distance a sensor can be placed from the collimator needs to be calculated. The sensor then needs to be within this distance to ensure that all rays from the collimator will be collected. The maximum working distance was calculated [5] for the collimator according to the following equation and Table 1:

$$R_{max} = \frac{f l_{col}}{d_{max}} (D_{col} - D)$$

Table 1. Parameters for Calculation Maximum Working Distance

| | Symbol | mm | in | Notes |
|--------------------------|--------|---------|--------|----------------------|
| Maximum Working Distance | Rmax | 6379.57 | 251.16 | |
| Focal Length Collimator | flcol | 1776.00 | 69.92 | ... +/-2 mm |
| Diameter of Collimator | Dcol | 304.80 | 12.00 | |
| Maximum Target Size | dmax | 49.50 | 1.95 | 35 mm x 35 mm target |
| System Aperture Diameter | D | 127.00 | 5.00 | SR5000W |

R_{\max} in a reflective collimator is measured from the primary output mirror; whereas, in a refractive collimator, it is from the last optic.

The atmospheric losses can be assumed identical by measuring in the same laboratory conditions (temperature, humidity, air currents, and so forth) and providing a setup that will have the distance [6] to the external and collimated blackbodies equal. Matching the distances also simplifies later calculations. A distance of 30 inches (") from the collimator aperture was selected to achieve a total of 158" (approximately 4 meters (m)) to allow the spectrometer to view an in-focus external blackbody. This distance still allowed flood-filling of the spectrometer FOV and matching the total atmospheric path length for the collimator, as shown in Figure 3. Another key setup parameter was ensuring that the SR5000 focus was set to infinity for the collimator and 4 m for the external blackbody.

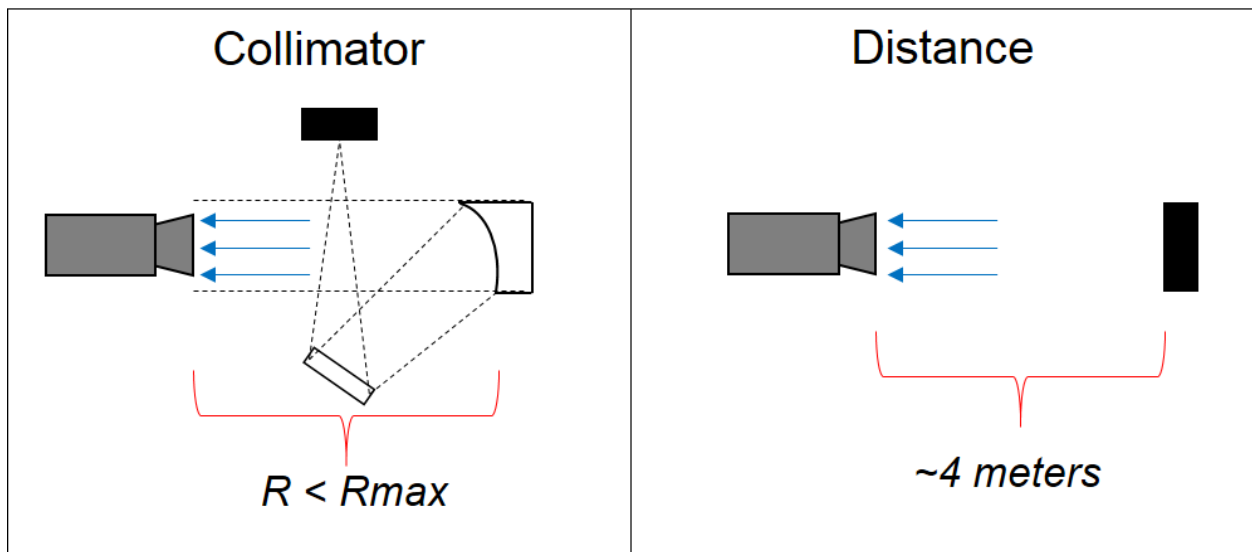


Figure 3. Setup of SR5000 With Atmospheric Paths Equal at 4 m

After the geometry was calculated, all instruments were setup and double checked for proper distances. Two blackbodies were used for ease of setup but selected very carefully. They were identical in vendor, make, model, size, and coating and had been fabricated within a year of each other. By doing this, the spectral emissivity contribution from the blackbodies were also canceled from the calculation. Both blackbodies were turned on, set to 100 °C, and allowed to stabilize, as shown in Figure 4. The temperature was chosen to allow plenty of signal to noise in the Long-Wave Infrared (LWIR) while still not saturating the instrument. The SR5000 was physically configured for spectral measurements, turned on, set to spectral mode, detector cooled with Liquid Nitrogen (LN2), and the entire system was allowed to stabilize for 20 minutes. Thirty spectrum scans at 16 seconds per scan were collected for both setups. The scans were then averaged to provide a single voltage versus wavelength data set.

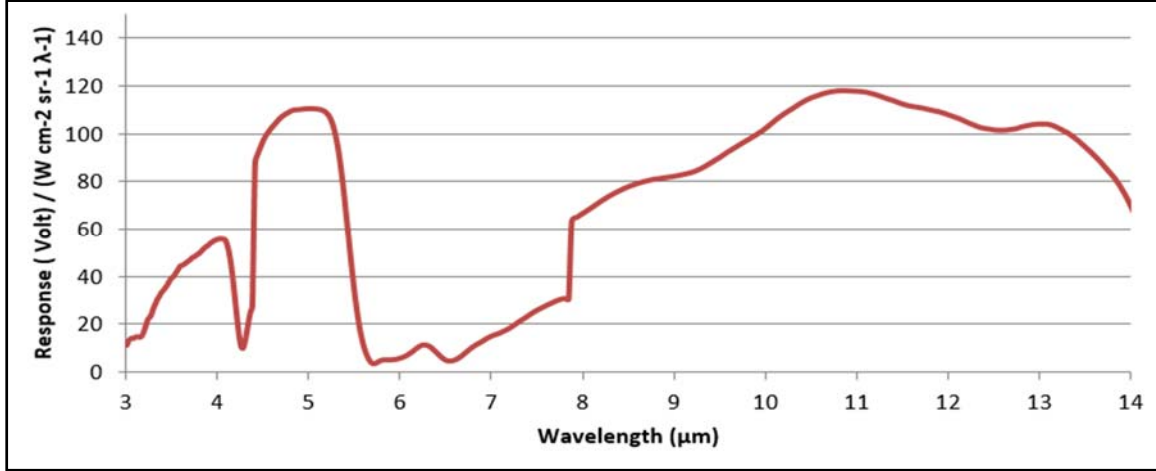


Figure 4. Response Function at 100 °C

The SR5000 has an internal reference blackbody that is sampled for every wavelength scan. It is used to compensate for any temperature drift inside the instrument housing. The following form of the Planck equation was used to calculate the radiance for both this internal blackbody and the external source:

$$L = \int_{\lambda_{min}}^{\lambda_{max}} \varepsilon(\lambda) \frac{C_1}{\lambda^5 \left(e^{\frac{C_2}{\lambda T}} - 1 \right)} d\lambda$$

where, $C_1 = 1.1909 \times 10^4$ (W μm^4 sr⁻¹ cm⁻²); $C_2 = 1.4388 \times 10^4$ (μm^4 K); and $\varepsilon(\lambda)$ = emissivity of the blackbody. The internal blackbody is assumed to be 1, while the external and collimator blackbodies are given as 0.97 ± 0.02 . This value will be measured at a later date, but for these calculations, it is not important since both the collimator and external blackbodies are identical models. The system response function is then calculated as follows:

$$R_{ext}(\lambda) = \frac{S_{ext}(\lambda)}{L(\lambda, T_{ext}) - L(\lambda, T_{IBB})}$$

where, S is the measured signal of known temperature filling the Field of View (FOV), R has units of volt per radiance, and L is the radiance of the external and internal blackbodies, respectively.

Since the external blackbody was used to calculate the response function, a correction factor must be used to account for the two different focus set points of 4 m for the external blackbody and infinity for the collimator, as shown in Figure 5. This is due to slight F Number (F/#) changes (for example, the primary mirror in the SR5000 changes position and therefore focal length), as shown in Figure 6:

$$R(\lambda) = \frac{R_{ext}(\lambda)}{((d - fl)/d)^2}$$

where, d is the calibration distance in meters; f is the focal length in meters (0.25 m for this SR5000W); R_{ext} is the measured response of the external blackbody at distance d .

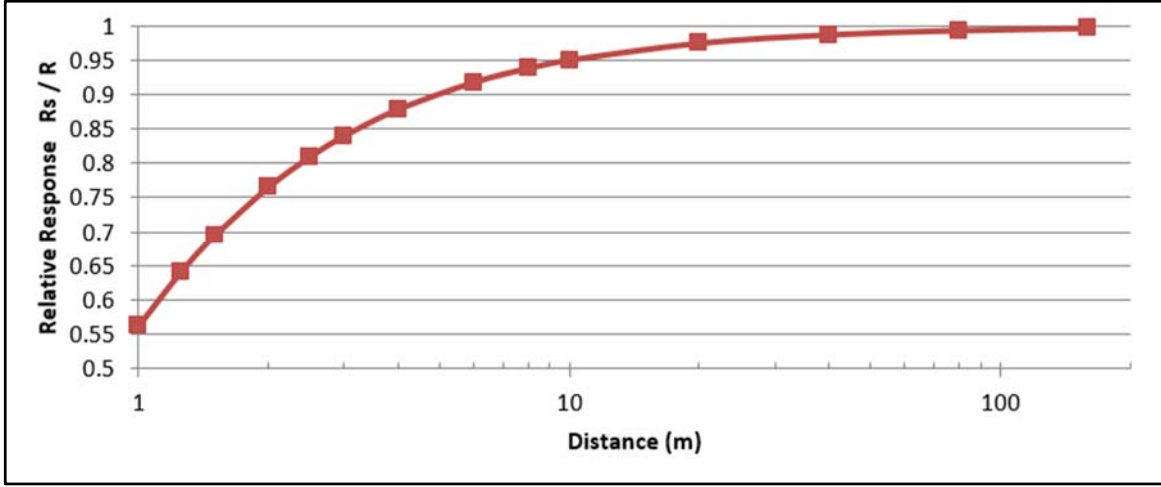


Figure 5. Response at Distance Divided by Response at Infinity

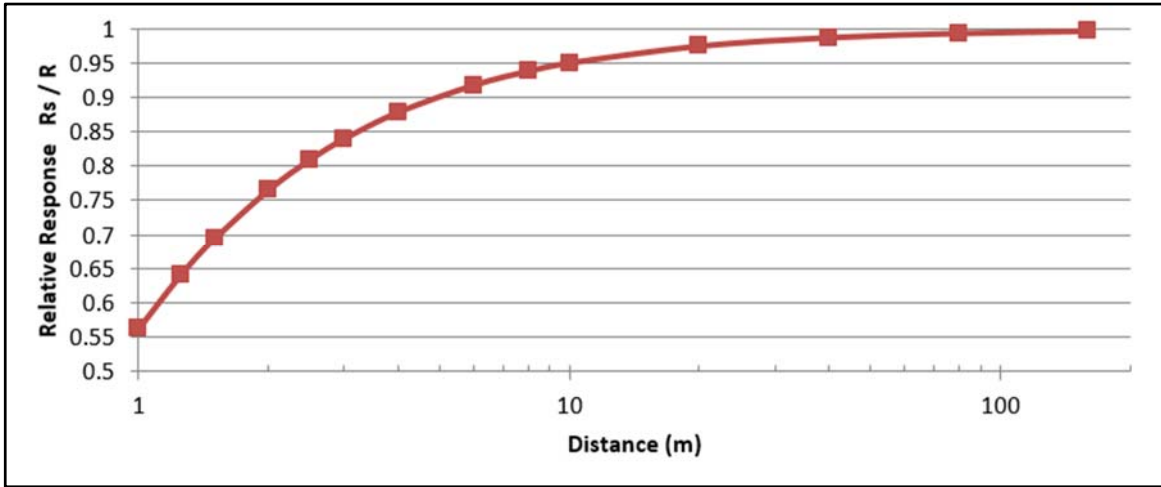


Figure 6. Change in F/# Due to Change in Focus for Each Distance

Once the response is known, the collimator source spectral radiance W can be calculated:

$$W(\lambda) = \frac{S_{col}(\lambda)/R(\lambda) + L(\lambda, T_{IBB}) - L(\lambda, T_{air})(1 - \tau(\lambda, T_{air}))}{\tau(\lambda, T_{air})}$$

where, S is the collimator signal; R is the previously calculated system response function with F/# correction; tau (τ) is the transmission of the atmosphere; and L is the spectral radiance of the internal blackbody and the air, respectively. Since the test setup was carefully chosen, the atmospheric effects were eliminated and equation simplified:

$$(\lambda) = \frac{S_{col}(\lambda)}{R(\lambda)} + L(\lambda, T_{IBB})$$

IV. RESULTS AND ANALYSIS

Plotting the spectral radiance, as shown in Figures 7 and 8, of the external blackbody versus the blackbody in the collimator clearly shows the losses coming from the mirrors.

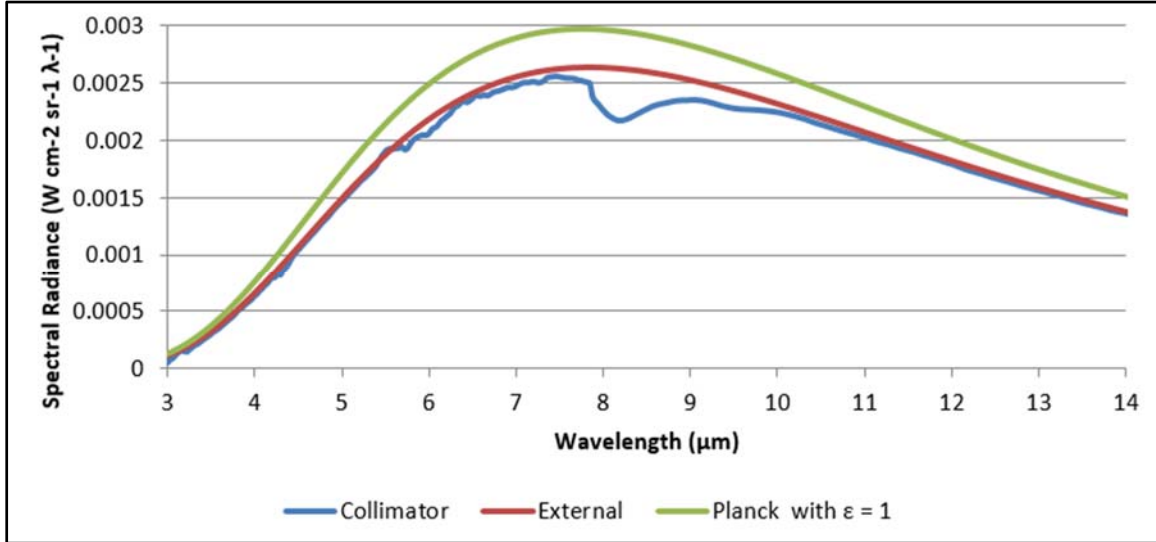


Figure 7. Radiance at 100 °C

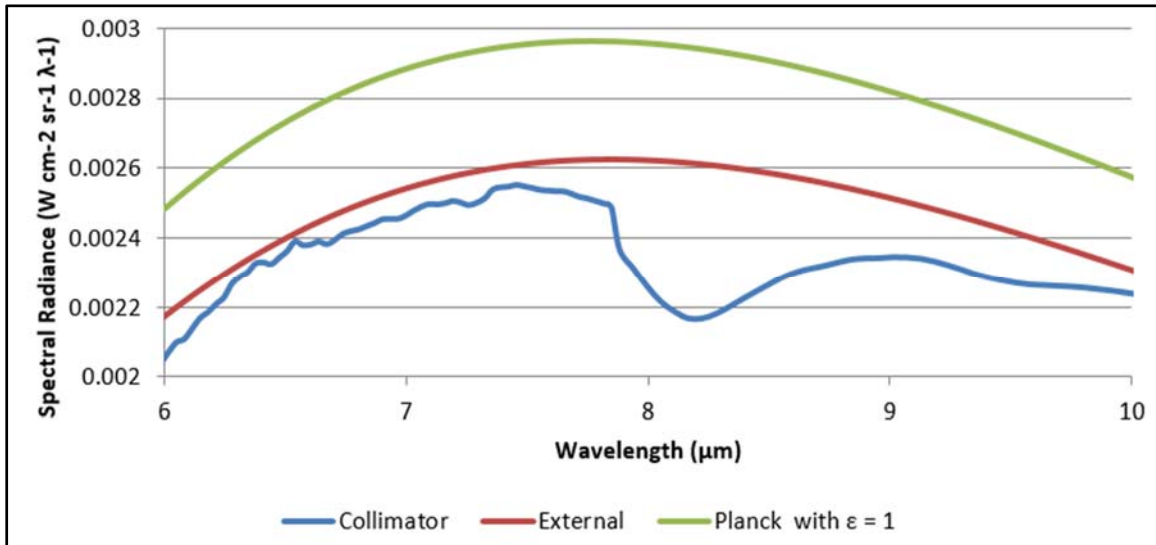


Figure 8. Radiance at 100 °C Zoomed-In

Larger than expected transmission losses are shown across the entire waveband, especially at 8 μm . Figure 9 compares the measured actual transmission to the vendor-supplied calculated transmission. The Mid-Wave (MW) band will require additional investigation, but it is suspected that the signal (for example, blackbody temperature) would need to be elevated to increase the signal to noise.

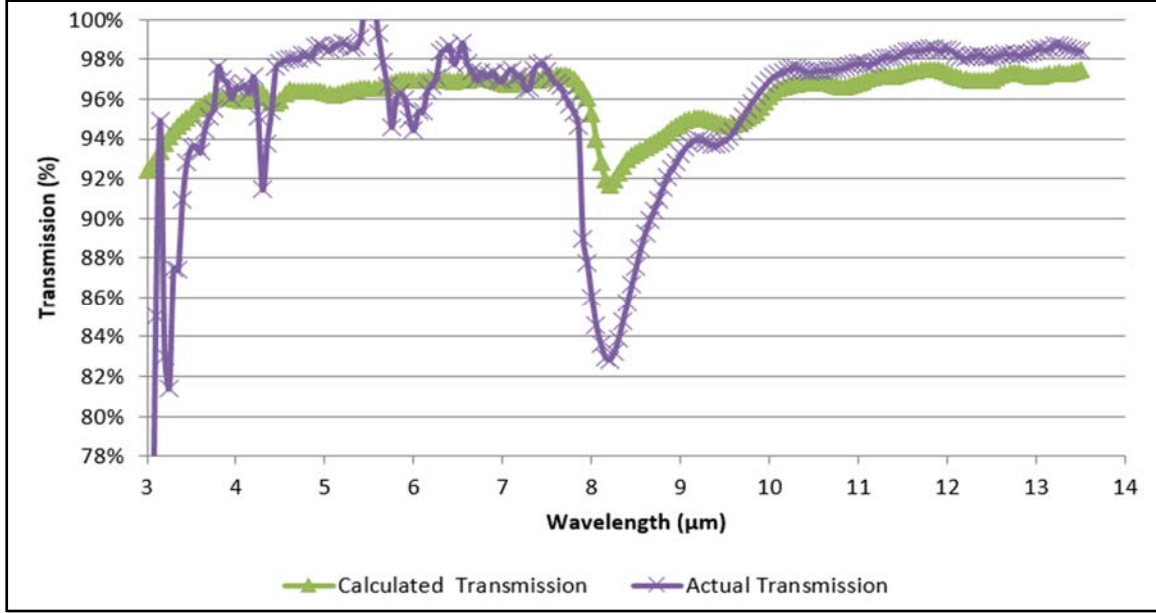


Figure 9. Reported Versus Measured Collimator Transmission

The transmission loss of example sub-bands can be calculated now that the losses for the collimator have been characterized. Table 2 shows large errors in transmission can be introduced if the assumed 97% (over 3 to 15 μm) is used instead of either the calculated or measured. Depending on the sub-band selected, large errors can still occur even if using the calculated transmission. The ideal solution can be obtained by using the measured, actual collimator transmission.

Table 2. Average Transmission of Several Sub-Bands

| Sub-Bands (μm) | Calculated Transmission | Actual Transmission | Additional Loss Assumed (97%)—Actual | Additional Loss Calculated—Actual |
|-----------------------------|-------------------------|---------------------|--------------------------------------|-----------------------------------|
| 7-13.5 | 96.2% | 95.6% | - 1.4 % | - 0.6% |
| 8-10 | 94.2% | 91.1% | - 5.9% | - 3.1% |
| 8-12 | 95.6% | 94.5% | - 2.5% | - 1.1% |
| 8-12.5 | 95.8% | 94.9% | - 2.1% | - 0.9% |
| 3.0-5.0 | 95.5% | 92.8% | - 4.2% | - 2.8% |
| 3.35-4.15 | 95.7% | 94.7% | - 2.3% | - 1.0% |

V. CONCLUSION AND FUTURE WORK

The collimator transmission losses can be measured by carefully selecting the test setup to eliminate atmospheric transmission and spectral emissivity differences of the blackbodies. This method detailed the equations necessary to calibrate and compensate for focus changes due to changing F/# of the SR5000 spectroradiometer. Measuring and using the actual collimator transmission losses resolved the discrepancies seen when comparing data through a collimator versus a standalone large surface blackbody. Further work will be done to secure a high emissivity cavity blackbody with a National Institute of Standards and Technology (NIST) measured spectral emissivity. This will then be used as the calibration source to completely characterize a collimator—mirrors spectral transmission and blackbody spectral emissivity.

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LIST OF ABBREVIATIONS, ACRONYMS, AND SYMBOLS

| | |
|------|--|
| % | percent |
| ” | inch |
| ~ | approximately |
| < | less than |
| > | greater than |
| ° | degree |
| μm | micrometers |
| C | Celsius |
| Cam | Camera |
| cm | centimeter |
| F/# | F-Number |
| fl | Focal Length |
| FOV | Field of View |
| IBB | Internal Blackbody |
| in | inch |
| InSb | indium antimonide |
| IR | Infrared |
| K | Kelvin |
| L | Radiance |
| LN2 | Liquid Nitrogen |
| LWIR | Long-Wave Infrared |
| m | meter |
| MCT | mercury cadmium telluride |
| mm | millimeter |
| MW | Mid-Wave |
| MWIR | Mid-Wave Infrared |
| NIST | National Institute of Standards and Technology |

LIST OF ABBREVIATIONS, ACRONYMS, AND SYMBOLS (CONCLUDED)

| | |
|---------------|---------------|
| SiO | silicon oxide |
| sr | steradian |
| T | Temperature |
| W | watt |
| x | times |
| ε | emissivity |
| λ | wavelength |
| τ | transmission |